# Quantum dynamics of a particle constrained to lie on a surface

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<sup>&</sup>lt;sup>1</sup>Work done under support of FAPESP.

## In this talk we will

- Recall the Schrödinger equation on  $\mathbb{R}^3$ .
- ▶ Present a Schrödinger-type equation on a surface of ℝ<sup>3</sup>.
- Relate the solution to the latter to the solution to the former.

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## Outline

- 1. Introduction
- 2. Main result
- 3. Sketch of proof

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Quantum mechanics for a particle in  $\mathbb{R}^3$ 

#### Wave function

$$\psi(x,t)\in\mathbb{C},\qquad x\in\mathbb{R}^3,\qquad t\in\mathbb{R}.$$

Square-integrable:

$$\psi(\cdot, t) \in L^2(\mathbb{R}^3, dx).$$

Normalized:

$$\int_{\mathbb{R}^3} |\psi(x,t)|^2 \, dx = 1.$$

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Quantum dynamics of a particle in  $\mathbb{R}^3$ 

Schrödinger equation

$$\begin{cases} i\partial_t \psi = H\psi \\ \psi|_{t=0} = \psi_0 \end{cases} \iff \psi(\cdot, t) = e^{-itH}\psi_0.$$

• Hamiltonian with  $A(x) \in \mathbb{R}^3$  and  $V(x) \in \mathbb{R}$ :

$$\begin{aligned} H\psi &= \left[ (i\nabla + A)^2 + V \right] \psi \\ &= i \operatorname{div}(i \operatorname{grad}(\psi) + A\psi) + \langle A, i \operatorname{grad}(\psi) + A\psi \rangle + V\psi. \end{aligned}$$

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# Calculus in a manifold $S^k$

Coordinates:

$$x_1(p),\ldots,x_k(p)$$
 for  $p\in S^k$ .

Metric:

$$[g_{ij}], \qquad [g^{ij}] = [g_{ij}]^{-1}, \qquad g = \det[g_{ij}].$$

Gradient:

$$(\operatorname{grad} f)^i = g^{ij} \frac{\partial}{\partial x_j} f.$$

Divergent:

div 
$$Y = \frac{1}{\sqrt{g}} \frac{\partial}{\partial x_j} (\sqrt{g} Y^j).$$

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What if the particle is constrained to lie on  $\Sigma$ ?

Suppose

 $x \in \Sigma$ , surface  $\Sigma \subset \mathbb{R}^3$ .

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What is the equation for  $\psi(x, t)$  on  $\Sigma$  ?

## What if the particle is constrained to lie on $\Sigma$ ?

Suppose

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What is the equation for  $\psi(x, t)$  on  $\Sigma$  ?

Natural candidate:

Replace  $\mathbb{R}^3$  and  $\langle \cdot , \cdot \rangle$  by  $\Sigma$  and  $\langle \cdot , \cdot \rangle|_{\Sigma}$ . On  $\Sigma$  we have:

 $L^2(\Sigma, dvol), div, grad, A, V, H.$ 

Equation:

$$\begin{cases} i\partial_t \psi = H\psi \\ \psi|_{t=0} = \psi_{0,\Sigma} \end{cases} \iff \psi(\cdot, t) = e^{-itH}\psi_{0,\Sigma}.$$

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Is this the right equation on  $\Sigma$ ?

 The solution of this equation on Σ must agree with Schrödinger on R<sup>3</sup>.

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Is this the right equation on  $\Sigma$ ?

- The solution of this equation on Σ must agree with Schrödinger on R<sup>3</sup>.
- How do we check?
- ▶ In ℝ<sup>3</sup>, consider

$$H_{\lambda} = H + \lambda^4 W$$
 on  $L^2(\mathbb{R}^3)$ 

with

$$\text{large } \lambda \in \mathbb{R}, \qquad \mathcal{W}|_{\Sigma} = 0, \qquad \mathcal{W}|_{\mathbb{R}^3 \setminus \Sigma} > 0.$$

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The potential  $\lambda^4 W$  traps the particle near  $\Sigma$ .

## We want to compare solutions

Is it true that



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where  $\psi_{0,\lambda}$  is supported near  $\Sigma$ ?

### We want to compare solutions

Is it true that



where  $\psi_{0,\lambda}$  is supported near  $\Sigma$ ?

No. But if replace H on L<sup>2</sup>(Σ) by

$$H_{\Sigma} = H + K$$
 on  $L^2(\Sigma)$ ,

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then yes.

### We want to compare solutions

Is it true that



where ψ<sub>0,λ</sub> is supported near Σ? ► No. But if replace *H* on L<sup>2</sup>(Σ) by

$$H_{\Sigma} = H + K$$
 on  $L^2(\Sigma)$ ,

then yes.

 $K = s - h^2 \equiv$  geometric potential,  $s \equiv$  Gaussian curvature,  $h \equiv$  mean curvature (not intrinsic).

s, h, K are functions of  $\sigma \in \Sigma$ .

## Examples

► Torus:

$$s = \frac{1}{r^2}, \qquad h = \frac{1}{r}, \qquad K = 0.$$

 $K \neq 0.$ 

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## Motivation

Effective evolution equations:

One-particle system embedded in a larger system.

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- Lab applications (physics literature).
- Comparison to classical dynamics.

# Previous work (among others)

## Physics

- ▶ da Costa, Phys. Rev. A, (1981). Others...
- ► Ferrari and Cuoghi, Phys. Rev. Lett. (2009).

Math-phys

► Froese and Herbst, Comm. Math. Phys. (2001).

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- ► Dell'Antonio and Tenuta (2004) (semiclassics).
- ▶ Wachsmuth and Teufel (2009) (adiabatic).

#### Our contribution

- ► More general *W*.
- Magnetic potential A.
- $H_{\Sigma}$  doesn't depend on  $A_{\perp \Sigma}$ .

Continuing... We will use

Normal bundle

$$\begin{split} \mathsf{N}\Sigma &= \{(\sigma, n) \mid \sigma \in \Sigma, \ n \in \mathsf{N}_{\sigma}\Sigma\},\\ \mathsf{N}\Sigma_{\delta} &= \{(\sigma, n) \mid |n| < \delta\}. \end{split}$$

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#### Change of variables

$$E: N\Sigma \to \mathbb{R}^3,$$
  
$$E(\sigma, n) = \sigma + n.$$

 $U_{\delta} \equiv \text{tubular neighborhood of } \Sigma \text{ in } \mathbb{R}^3,$  $E: N\Sigma_{\delta} \to U_{\delta}$  is a diffeomorphism.

## Rewriting the problem

We want to study

$$e^{-itH_{\lambda}}\psi_{0,\lambda}$$
 on  $L^2(\mathbb{R}^3)$  near  $\Sigma$  (Euclidean metric).

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We do

Restrict to  $L_0^2(U_{\delta})$ Error  $O(\lambda^{-1})$ Change variables to  $L^2(N\Sigma_{\delta})$ No errorExtend to  $L^2(N\Sigma)$ Error  $O(\lambda^{-1})$ 

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We end up with

$$e^{-itH_{\lambda}}\psi_{0,\lambda}$$
 on  $L^2(N\Sigma)$  (metric  $g_{N\Sigma}$ ).

Decomposition of velocities

Tangent space of  $N\Sigma$  at  $(\sigma, n)$ 

 $T_{\sigma,n}N\Sigma$ .

Horizontal and vertical components of  $T_{\sigma,n}N\Sigma$ 

$$T_{\sigma,n}N\Sigma \simeq (T_{\sigma,n}N\Sigma)_H \oplus (T_{\sigma,n}N\Sigma)_V,$$
  
(X,Y) =  $P_H(X,Y) + P_V(X,Y).$ 

 $\langle (X_1, Y_1), (X_2, Y_2) \rangle_{N\Sigma_{\delta}} = \langle X_1 + Y_1, X_2 + Y_2 \rangle_{\mathbb{R}^3}.$ 

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## Gauge transformation

Function  $\gamma$  on  $N\Sigma$ . Unitary transformation:

$$egin{aligned} S_\gamma &= e^{i\gamma} \quad ext{on} \quad L^2(N\Sigma),\ S_\gamma^* \mathcal{H}_\lambda S_\gamma,\ A &\leadsto A - ext{grad}(\gamma). \end{aligned}$$

#### Proposition

$$P_H(A - \operatorname{grad}(\gamma))|_{\Sigma} = P_H A,$$
  
 $P_V(A - \operatorname{grad}(\gamma)) = 0.$ 

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#### Proposition

$$\begin{aligned} & P_H(A - \operatorname{grad}(\gamma))|_{\Sigma} = P_H A, \\ & P_V(A - \operatorname{grad}(\gamma)) = 0. \end{aligned}$$

Proof:

$$\gamma(x,y) = \int_0^y A_3(x,s) \, ds$$

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does the job in each chart. Check that works globally.

## Dilation in the normal directions

Unitary transformation:

$$egin{aligned} &U_\lambda: L^2(N\Sigma, \operatorname{dvol}_{N\Sigma}) o L^2(N\Sigma, \operatorname{dvol})\ &(U_\lambda \psi)(\sigma, n) = \sqrt{\lambda} \, m(\sigma, n) \psi(\sigma, \lambda n) \end{aligned}$$

Ignore dvol<sub>N $\Sigma$ </sub> and  $m(\sigma, n) > 0$  for the moment.

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Initial data

 $e^{-itH_{\lambda}}\psi_{0,\lambda}$  on  $L^2(N\Sigma, dvol), \qquad \psi_{0,\lambda} = S_{\gamma}U_{\lambda}\psi_0.$  $\psi_{0,\lambda}$  is squeezed towards  $\Sigma$ .

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Plot of  $\psi(\sigma, \cdot)$  (blue) and  $(U_{\lambda}\psi)(\sigma, \cdot)$  (red)



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## Putting all together

#### We arrive at

$$e^{-itL_{\lambda}}\psi_0$$
 on  $L^2(N\Sigma, dvol_{N\Sigma})$ 

with

$$L_{\lambda} = U_{\lambda}^* S_{\gamma}^* H_{\lambda} S_{\gamma} U_{\lambda}.$$

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## Large $\lambda$ expansion

#### Local coordinates

$$\Sigma : \quad x(\sigma)$$

$$N\Sigma : \quad \begin{cases} x(\sigma, n) = x(\sigma) \\ y(\sigma, n) = \langle \nu(\sigma), n \rangle \end{cases}$$

$$TN\Sigma : \quad \partial/\partial x, \, \partial/\partial y$$

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Expand the metric in  $y\lambda^{-1}$ 

$$g_{N\Sigma_{\delta}}(x,y\lambda^{-1}) = \begin{pmatrix} G_{\Sigma}(x)(I-y\lambda^{-1}M(x))_{2\times 2} & 0\\ 0 & 1 \end{pmatrix}_{3\times 3}$$

## Hamiltonians

#### Hypothesis and notation

- A, V, W smooth functions of  $(\sigma, n)$ .
- Set  $\mathbb{A}(\sigma) = P_H A(\sigma, 0)$  and  $\mathbb{V}(\sigma) = V(\sigma, 0)$ .

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## Hamiltonians

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Expansion for  $L_{\lambda}$ 

$$L_{\lambda} = H_{\Sigma} + \lambda^2 H_{O,\lambda} + O(\lambda^{-1}).$$

Hamiltonian on  $\Sigma$ :

 $H_{\Sigma}\psi = i\operatorname{div}_{\Sigma}(i\operatorname{grad}_{\Sigma}(\psi) + \mathbb{A}\psi) + \langle \mathbb{A}, i\operatorname{grad}_{\Sigma}(\psi) + \mathbb{A}\psi\rangle_{\Sigma} + (\mathbb{V} + K)\psi.$ 

Oscillator in the normal directions:

$$H_{O,\lambda}\psi = -\Delta_n\psi + (\langle n, Bn \rangle_{\mathbb{R}^3} + \underbrace{O(\lambda^{-1})}_{\psi})\psi.$$

will not depend on  $\sigma$ 

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## Main result

## Theorem (J. Math. Phys. (2014))

W has local minimum at Σ.

• 
$$\partial_{\sigma}W(\sigma,n) = O(|n|^5)$$
 as  $|n| \to 0$ .

- $W(\sigma, n) \geq c|n|^2$ .
- *g*<sub>NΣ</sub> complete metric.

Then for any  $\psi_0$ , T > 0, and large  $\lambda$ :

$$\sup_{t\in[0,T]} \left\| e^{-itL_{\lambda}}\psi_0 - e^{-it(H_{\Sigma}+\lambda^2 H_{O,\lambda})}\psi_0 \right\| \leq \frac{C}{\sqrt{\lambda}}.$$

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#### Remarks

- Neither orbits converge. Only their difference.
- Implies a statement on  $\mathbb{R}^3$ .

## Interpretation

Since

$$L^2(N\Sigma, \operatorname{dvol}_{N\Sigma}) = L^2(\Sigma, \operatorname{dvol}) \otimes L^2(\mathbb{R}, dy),$$
  
 $[H_{\Sigma}, H_{O,\lambda}] = 0,$ 

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we have

$$H_{\Sigma} = h_{\Sigma} \otimes I,$$
  
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we have

$$H_{\Sigma} = h_{\Sigma} \otimes I,$$
  
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and

$$\begin{split} \exp(-it(H_{\Sigma}+\lambda^2 H_{O,\lambda})) &= \exp(-itH_{\Sigma})\exp(-it\lambda^2 H_{O,\lambda}) \\ &= \exp(-ith_{\Sigma})\otimes\exp(-it\lambda^2 h_{O,\lambda}). \end{split}$$

Motion on  $\Sigma$  superposed by normal oscillations.

## Sketch of proof

Energy cutoff:

 $E_{<}$  equals 1 if  $\langle \psi, L_{\lambda}\psi \rangle < \mu\lambda^{2}$  and 0 otherwise.

► |*n*| cutoff:

 $N_{<}$  equals 1 if  $|n| < \delta \lambda$  and 0 otherwise.

Partition of unity for Σ:

 $\{\chi_j\}.$ 

Resolution of identity:

$$1 = 1 \cdot 1 \cdot 1$$
  
=  $\left[\sum_{j=1}^{m} \chi_j(\sigma)^2\right] \cdot \left[N_{<} + N_{\geq}\right] \cdot \left[E_{<} + E_{\geq}\right]$   
=  $\chi_1 \cdot N_{<} \cdot E_{<}$  + Reminder.

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## We want to estimate

Notation: 
$$L_{0,\lambda} = H_{\Sigma} + \lambda^2 H_{O,\lambda}$$
.  
 $\|e^{-itL_{\lambda}}\psi_0 - e^{-itL_{0,\lambda}}\psi_0\|^2$   
 $\leq \left|\int_0^t ds \frac{d}{ds} \langle e^{-isL_{\lambda}}\psi_0, 1 \cdot e^{-isL_{0,\lambda}}\psi_0 \rangle\right|$   
 $\leq Ct \sup_s \langle \chi_1 E_{<} e^{-isL_{\lambda}}\psi_0, \underbrace{[L_{0,\lambda}, N_{<}]}_{\sim \{\partial_y N_{<}\}D_y} + N_{<}(\underbrace{L_{\lambda} - L_{0,\lambda}}_{\sim D_x})e^{-isL_{0,\lambda}}\psi_0 \rangle$   
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 $+ O(\lambda^{-1})$ 

Using Cauchy-Schwarz, we reduce the problem to

$$\begin{split} \lambda^2 \| (\chi_1 N_{\leq} E_{\leq}) \left\{ \partial_y N_{\leq} \right\} D_y e^{-isL_{0,\lambda}} \psi_0 \| &\leq C \lambda^{-1} \quad (\text{Energy bounds}). \\ \| (\chi_1 N_{\leq} E_{\leq}) \underbrace{D_x e^{-isL_{0,\lambda}}}_{[D_x,W]} \psi_0 \| &\leq C \lambda^{-1/2} \quad (\text{Propagation bounds}). \end{split}$$

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#### Appendix

## Energy bounds

#### Schematically

We have

$$D_x^*D_x+V(x,y)+\mathcal{K}(x,y)+\lambda^2(D_y^*D_y+\lambda^2\mathcal{W}(x,y\lambda^{-1}))=L_\lambda\leq\lambda^2\mu.$$

Use conservation of expected value of energy and positivity. We obtain

$$(\lambda^{-1}D_x^*)^{\alpha}(\lambda^{-1}D_x)^{\alpha}+(D_y^*)^pD_y^p+(\lambda^2W(x,y\lambda^{-1}))^l\leq C(\lambda^{-2}L_\lambda)^{l+1},$$

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where  $|\alpha| + p \leq 2$  and  $l \geq 1$ .

## Propagation bounds

Let

$$f(t) \equiv$$
 expected value of  $D_x^* D_x$  on  $e^{-itL_\lambda} \psi_0$ .

Use:

Gronwall's inequality:

$$rac{d}{dt}f(t)\leq Cf(t), \qquad f(0)\leq C \quad \Longrightarrow \quad \sup_t f(t)\leq C.$$

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Energy bounds.

Thank you for your attention.