Quantitative derivation of the Gross-Pitaevskii equation

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This talk is about

Mathematics of many-body quantum mechanics.

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- Dynamics of Bose-Einstein condensates.
- Effective description.
- How the Gross-Pitaevskii PDE emerges.

Wave function for N bosons

N-particle wave function:

$$\psi_N(x_1,\ldots,x_N,t)\in\mathbb{C},\qquad x_1,\ldots,x_N\in\mathbb{R}^3,\qquad t\in\mathbb{R}.$$

Square-integrable and normalized:

$$\psi_N(\cdot,t)\in L^2(\mathbb{R}^{3N})$$
 and $\int_{\mathbb{R}^{3N}}|\psi_N(\cdot,t)|^2=1.$

• ψ_N is symmetric in each pair of variables x_1, \ldots, x_N .

Density operator

$$|\psi_N\rangle\langle\psi_N|$$
 on $L^2(\mathbb{R}^{3N})\simeq L^2(\mathbb{R}^3)\otimes\cdots\otimes L^2(\mathbb{R}^3).$

Bose-Einstein condensate

In experiments (since 1995)

Trapped cold ($T \sim 10^{-9} K$) dilute gas of $N \sim 10^3$ bosons.

Heuristically

$$\psi_N(x_1,\ldots,x_N,t_0) \simeq \prod_{j=1}^N \varphi(x_j) \quad \text{where} \quad \varphi \in L^2(\mathbb{R}^3).$$

 $|\psi_N \rangle \langle \psi_N | \simeq |\varphi \rangle \langle \varphi | \otimes \cdots \otimes |\varphi \rangle \langle \varphi |.$

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Condensate states

One-particle reduced density operator

$$|\psi_N\rangle\langle\psi_N|^{(1)} = \operatorname{Trace}_{2\to N}|\psi_N\rangle\langle\psi_N|.$$

(Integrate out N-1 variables of the integral kernel of $|\psi_N\rangle\langle\psi_N|$.)

 $|\psi_N\rangle\langle\psi_N|^{(1)}$ plays the role of one-particle wave-function.

Condensate condition

$${\sf Tr} \ \Big| \ |\psi_{\sf N}
angle \langle \psi_{\sf N}|^{(1)} - |arphi
angle \langle arphi| \ \Big| \leq rac{{\sf C}}{{\sf N}} o 0 \quad {\sf as} \quad {\sf N} o \infty.$$

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Model (which is realistic)

Quantum Hamiltonian in the Gross-Pitaevskii regime

$$H_N^{\mathrm{trap}} = \sum_{j=1}^N \left(-\Delta_{x_j} + V_{\mathrm{trap}}(x_j) \right) + \frac{1}{N} \sum_{i < j}^N N^3 V(N(x_i - x_j)),$$

 $V_{ ext{trap}}(y) = |y|^2$ and $V \ge 0$ with compact support.

Very heuristically

$$rac{1}{N}N^{3}V(N\cdot)\simrac{1}{N}\delta(\cdot)$$
 for large N

models rare but strong collisions.

Time evolution of condensates

Initial condition

 $\psi_N|_{t=0} = \theta_N = \text{condensate state with correlations (not a product)}$

We construct initial data Θ in Fock space:

$$\Theta = W(\sqrt{N}\varphi)T(k)\Omega = \theta_0 \oplus \theta_1 \oplus \cdots \oplus \theta_N \oplus \cdots \oplus \bigoplus_{n\geq 0} L^2_{sym}(\mathbb{R}^{3n})$$

 $\Omega = \text{finite particle state (e.g. vacuum)}$ T(k) = Bogoliubov transformationk(x, y) = integral kernel which models correlations $W(\sqrt{N}\varphi) = \text{Weyl operator}$ $\varphi(x) = \text{one particle state}$ $\Theta = \text{modified coherent state}$

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Schrödinger equation on Fock space

Condensate state reached; traps are turned off

$$H_N = H_N^{
m trap}$$
 with $V_{
m trap} \equiv 0.$

Hamiltonian on Fock space

$$\mathcal{H}=H_0\oplus H_1\oplus\cdots\oplus H_N\oplus\cdots$$

Time evolution is observed

$$\begin{cases} i\partial_t \Psi = \mathcal{H}\Psi \\ \Psi|_{t=0} = W(\sqrt{N}\varphi)T(k)\Omega \end{cases} \quad \text{as} \quad N \to \infty.$$

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Theorem [Benedikter, deO, Schlein, CPAM 2014]

Reasonable hypothesis on $V \ge 0$, φ , Ω . Consider the solution

$$\Psi = e^{-i\mathcal{H}t}W(\sqrt{N}\varphi)T(k)\Omega.$$

Let

$$\Gamma_{N,t}^{(1)} =$$
 one-particle reduced operator of Ψ .

Then

$$\mathsf{Tr} \, \Big| \, \mathsf{\Gamma}_{\mathsf{N},t}^{(1)} - |arphi
angle \langle arphi| \, \Big| \leq C \exp(C \exp(C|t|)) rac{1}{\sqrt{N}}$$

for all t and N, where φ_t solves (time-dep. Gross-Pitaevskii eqn.)

$$i\partial_t \varphi_t = -\Delta \varphi_t + 8\pi a_0 |\varphi_t|^2 \varphi_t$$
 with $\varphi_t|_{t=0} = \varphi_t$

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 $a_0 > 0$ (scattering length of V).

Remarks

Based on

► Hepp '74, Ginibre–Velo '79, Rodnianski–Schlein '09,...

Previous results

 Spohn '80, Erdös–Schlein–Yau '06, Pickl '10,... (no rate of convergence)

Other results

Adami–Golse–Teta '07, Grillakis–Machedon–Margetis '10,...

Large bibliography...

Look at arXiv:1208.0373 and Schlein's notes arXiv:1210.1603.

Thank you for your attention!